A Study of Shielding Properties of X-ray and Gamma in Barium Compounds

L. Seenappa¹, H.C. Manjunatha¹*, B.M. Chandrika², Hanumantharayappa Chikka³

¹Department of Physics, Government College for women, Karnataka, India; ²PC Extension, St.Annes School, Karnataka, India; ³Vivekananda Degree College, Karnataka, India

Background: Ionizing radiation is known to be harmful to human health. The shielding of ionizing radiation depends on the attenuation which can be achieved by three main rules, i.e. time, distance and absorbing material.

Materials and Methods: The mass attenuation coefficient, linear attenuation coefficient, Half Value Layer (HVL) and Tenth Value Layer (TVL) of X-rays (32 keV, 74 keV) and gamma rays (662 keV) are measured in Barium compounds.

Results and Discussion: The measured values agree well with the theory. The effective atomic numbers (Zeff) and electron density (Ne) of Barium compounds have been computed in the wide energy region 1 keV to 100 GeV using an accurate database of photon-interaction cross sections and the WinXCom program.

Conclusion: The mass attenuation coefficient and linear attenuation coefficient for BaCO₃ is higher than the BaCl₂, Ba(NO₃)₂ and BaSO₄. HVL, TVL and mean free path are lower for BaCO₃ than the BaCl₂, Ba(NO₃)₂ and BaSO₄. Among the studied barium compounds, BaCO₃ is best material for x-ray and gamma shielding.

Keywords: X-ray, Gamma, Barium compounds

Introduction

Ionizing radiation is known to be harmful to human health. The shielding of ionizing radiation depends on the attenuation which can be achieved by three main rules, i.e. time, distance and absorbing material. Shielding is the most effective technique and it is required against X-rays and gamma rays, which are very penetrating [1, 2]. Gamma radiation is best absorbed by dense materials consisting of heavy atoms such as lead and barium [3]. Barium is a dense alkaline earth metal in Group IIA- (alkaline earth metals) of the periodic table that occurs in nature as a divalent cation in combination with other elements. Two commonly found forms of barium are barium sulfate and barium carbonate, often found as underground ore deposits.

Shielding aprons manufactured from environmentally-friendly barium sulfate as a potential substitute for the lead aprons. Barium sulfate shielding can be processed easily to any type, with the same shielding ability as lead. Barium sulfate has long been utilized for in-vivo tests by employing the difference in density between soft tissues in radiology, and it is one of the most economic environment-friendly shielding materials [4]. Hubbell-Seltzer [5] tabulated the mass attenuation co-efficient for elements of
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In the present work, the mass attenuation coefficients and photon interaction cross sections in the energy range from 1 keV to 100 GeV are generated using WinXCom [7, 8]. The total molecular cross section $\sigma_m$ [milli barn], is computed from the following equation using the values of mass attenuation coefficients $[(\mu/\rho)]$,

$$\sigma_m(E) = \left( \frac{1}{N} \right) \left( \frac{\mu(E)}{\rho} \right) \sum_i n_i A_i$$

(1)

where $n_i$ is the number of atoms of $i^{th}$ element in a given molecule, $(\mu/\rho)c$ is the mass attenuation coefficient of bio-molecule, N is the Avogadro’s number, and $A_i$ is the atomic weight of element i. The effective (average) atomic cross section for a particular atom in the compound $\sigma_a$ [milli barn] is estimated using the equation,

$$\sigma_a = \sum_i n_i \sigma_m$$

(2)

The effective electronic cross section $\sigma_e$ [milli barn] is computed from mass attenuation coefficient $(\mu/\rho)$ of $i^{th}$ element in the given molecule using following equation,

$$\sigma_e = \left( \frac{1}{N} \right) \left( \frac{f_i A_i (\mu/\rho)}{Z_i} \right) = \frac{\sigma_a}{Z_{eff}}$$

(3)

where, $f_i$ is the fractional abundance (a mass fraction of the $i^{th}$ element in the molecule) and $Z_i$ is the atomic number of the $i^{th}$ element in a molecule. Finally the $Z_{eff}$ is estimated as

$$Z_{eff} = \frac{\sigma_a}{\sigma_e}$$

(4)

The effective electron density, $N_e$ expressed in terms of number of electrons per unit mass is closely related to the effective atomic number. For an element, the electron density is given by $N_e = N Z/A$ [electrons/g], This expression can be generalized for a compound,

$$N_e = \sum_i \frac{N}{n_i A_i} Z_{eff} \sum_i n_i$$

(5)

In the present work, $N_e$ and $Z_{eff}$ of Barium compounds are computed from Equations (4) and (5).

Materials and Methods

1. Present work

   1) Theoretical calculation of effective atomic number and electron density

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   In the present work, $N_e$ and $Z_{eff}$ of Barium compounds are computed from Equations (4) and (5).

2) Measurement of mass attenuation coefficient of gamma and other derived gamma attenuation parameters

   Transmission experiments with the narrow beam (good-geometry) setup were used for measuring the incident and transmitted intensities, and hence calculating the attenuation coefficient. The narrow geometry experimental setup used in the present measurement is as shown in Figure 1. We have used a NaI(Tl) crystal detector mounted on a photo-
multiplier tube housed in a lead chamber and a sophisticated PC based MCA for a detection purpose, gamma source such as $^{137}$Cs (661.6 keV) and Barium compounds such as BaCl$_2$, BaCO$_3$, Ba(NO$_3$)$_2$ and BaSO$_4$ as target samples. The sample was directly attached to the opening of the lead shield where source is placed. The integral intensities, $I_0$ and I of the beam before and after passing through the sample are measured for sufficient time. $(\mu/\rho)c$ of the sample is then estimated using the relation.

$$\left(\frac{\mu}{\rho}\right)_c = \left(\frac{I}{I_0}\right)\ln\left(\frac{I_o}{I}\right)$$  \hspace{1cm} (6)$$

Where, t and $\rho$ are the thickness and density of the sample respectively. Experimental values of $N_e$ and $Z_{eff}$ of Barium compounds are obtained by substituting the measured values of $(\mu/\rho)c$ in Equations (1)-(5). Some of the earlier workers measured the mass attenuation coefficient of photons using alternate to transmission experiments [9, 10]. The high-purity Germanium spectrometers having poor efficiency but high resolution, NaI scintillation detectors have a poor energy resolution but high efficiency [11]. For intensity measurements, efficiency of the detector should be high rather than resolution. Many previous [12-15] workers used the NaI scintillation detectors for the measurement of mass attenuation coefficients. Hence we have used the NaI scintillation detectors for the measurements.

**Results and Discussion**

The measured $^{137}$Cs gamma spectrum in different thickness of target samples such as BaCl$_2$, BaCO$_3$, Ba(NO$_3$)$_2$ and BaSO$_4$ are shown in the Figures 2-5. The experimental mass

![Fig. 2. The measured $^{137}$Cs spectrum in different thickness of BaCl$_2$ sample.](image1)

![Fig. 3. The measured $^{137}$Cs spectrum in different thickness of Ba(NO$_3$)$_2$ sample.](image2)

![Fig. 4. The measured $^{137}$Cs spectrum in different thickness of Ba(NO$_3$)$_2$ sample.](image3)

![Fig. 5. The measured $^{137}$Cs spectrum in different thickness of BaSO$_4$ sample.](image4)
attenuation coefficients are compared with the theoretical values and it is shown in the Table 1. This comparison shows that the measured values agree well with the theoretical values.

The total linear attenuation coefficient (μ) can be evaluated by multiplying density of compounds to mass attenuation coefficients.

\[ \mu = \left( \frac{\text{cm}^{-1}}{\text{g}} \right) \times \rho \]  

(7)

The total linear attenuation coefficient (μ) is used in the calculation of HVL. HVL is the thickness of a shield or an absorber that reduces the radiation level by a factor of 2 that is to half the initial level and is calculated by the following equation

\[ HVL = \frac{\ln 2}{\mu} = \frac{0.693}{\mu} \]  

(8)

The total linear attenuation coefficient (μ) is also used in the calculation of TVL. It is the thickness of a shield required for attenuating a radiation beam to 10% of its radiation level and is computed by

\[ TVL = \frac{\ln 10}{\mu} = \frac{2.303}{\mu} \]  

(9)

The average distance between two successive interactions is called the relaxation length (λ). It is also called the photon mean free path which is determined by the equation:

\[ \lambda = \frac{\int_0^\infty x \cdot \exp(-\mu x) \, dx}{\int_0^\infty \exp(-\mu x) \, dx} = \frac{1}{\mu} \]  

(10)

The measured linear attenuation coefficients (cm\(^{-1}\)), HVL (in cm), TVL (in cm) and mean free path (in cm) are compared with the theoretical values. This comparison is as shown in the Tables 2-5. The theoretical mass attenuation coefficients for the given compounds are computed for wide energy range 1 keV-100 GeV. These values are used in the calculation of linear attenuation coefficients (cm\(^{-1}\)), HVL (in cm), TVL (in cm) and mean free path (in cm) for wide energy range 1 keV-100 GeV. Figures 6-8 show that the variation of HVL, TVL and mean free path for different barium compounds with energy. From these figures, it is found that for all compounds the values of HVL, TVL and mean free path re-

<table>
<thead>
<tr>
<th>Table 1. Comparison of Measured Mass Attenuation Coefficients (cm(^2)/g) with Theoretical Values</th>
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<tbody>
<tr>
<td>Target compound</td>
</tr>
<tr>
<td>BaCl(_2)</td>
</tr>
<tr>
<td>BaCO(_3)</td>
</tr>
<tr>
<td>Ba(NO(_3))(_2)</td>
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<tr>
<td>BaSO(_4)</td>
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<tr>
<th>Table 2. Comparison of Measured Linear Attenuation Coefficients (cm(^{-1})) with Theoretical Values</th>
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<tr>
<td>Target compound</td>
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</tr>
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<td>BaSO(_4)</td>
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<tr>
<th>Table 3. Comparison of HVL (cm) with Theoretical Values</th>
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<tbody>
<tr>
<td>Target compound</td>
</tr>
<tr>
<td>BaCl(_2)</td>
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<tr>
<td>Ba(NO(_3))(_2)</td>
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<td>BaSO(_4)</td>
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Table 4. Comparison of TVL (cm) with Theoretical Values

<table>
<thead>
<tr>
<th>Target compound</th>
<th>32 keV</th>
<th>74 keV</th>
<th>661.7 keV</th>
</tr>
</thead>
<tbody>
<tr>
<td>BaCl₂</td>
<td>0.07E-02 ± 1.77E-03</td>
<td>0.77E-02</td>
<td>0.90E-01 ± 4.76E-03</td>
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<tr>
<td>BaCO₃</td>
<td>5.13E-02 ± 1.29E-03</td>
<td>5.64E-02</td>
<td>1.60E-01 ± 3.99E-03</td>
</tr>
<tr>
<td>BaNO₃</td>
<td>7.31E-02 ± 1.83E-03</td>
<td>7.96E-02</td>
<td>1.86E-01 ± 4.65E-03</td>
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<tr>
<td>BaSO₄</td>
<td>9.56E-02 ± 2.59E-03</td>
<td>1.03E-01</td>
<td>2.57E-01 ± 6.42E-03</td>
</tr>
</tbody>
</table>

Table 5. Comparison of Photon Mean Free Path (cm) with Theoretical Values

<table>
<thead>
<tr>
<th>Target compound</th>
<th>32 keV</th>
<th>74 keV</th>
<th>661.7 keV</th>
</tr>
</thead>
<tbody>
<tr>
<td>BaCl₂</td>
<td>3.07E-02 ± 7.68E-04</td>
<td>3.38E-02</td>
<td>8.27E-02 ± 2.07E-03</td>
</tr>
<tr>
<td>BaCO₃</td>
<td>2.23E-02 ± 5.58E-04</td>
<td>2.45E-02</td>
<td>6.93E-02 ± 1.73E-03</td>
</tr>
<tr>
<td>BaNO₃</td>
<td>3.18E-02 ± 7.94E-04</td>
<td>3.46E-02</td>
<td>8.08E-02 ± 2.02E-03</td>
</tr>
<tr>
<td>BaSO₄</td>
<td>4.15E-02 ± 1.04E-03</td>
<td>4.46E-02</td>
<td>1.12E-01 ± 2.79E-03</td>
</tr>
</tbody>
</table>

Fig. 6. The variation of Half Value Layer (HVL) with energy for different barium compounds.

Fig. 7. The variation of Tenth Value Layer (TVL) with energy for different barium compounds.

Fig. 8. The variation of Tenth Value Layer means free path or relaxation length with energy for different barium compounds.

remains constant from 1 keV to 10 keV. These parameters increase sharply up to few MeV and thereafter shows decreasing trend. From these figures it is observed that the HVL, TVL and mean free path increases up to the E_{pe} and then decreases [16]. Here E_{pe} is the energy value at which the photoelectric interaction coefficients matches with Compton interaction coefficients. The variation of these parameters with energy is due to dominance of photoelectric absorption in the lower end and dominance of pair production in the higher photon energy region. In the lower energy end photoelectric absorption is dominant, hence these values are minimum. As the energy of incident photon increases, Compton scattering overtakes the photoelectric absorption. It results multiple Compton scattering events which increases these values up to the E_{pe} and becomes maximum at E_{pe}. Thereafter
(above \(E_{\text{th}}\)), pair production starts dominating (absorption process) which reduces these parameters to minimum value.

The variation of \(Z_{\text{eff}}\) with photon energy for total photon interactions is as shown in the figure 9 and this variation is because of dominance of different photon interactions. \(Z_{\text{eff}}\) increases and becomes maximum and decreases sharply in the energy region 0.002-0.035 MeV. The variation of \(Z_{\text{eff}}\) with energy is constant for 0.5 MeV to 2 MeV, from 2 MeV to 70 MeV increases and thereafter remains constant. The effective electron density shows similar trend as that of \(Z_{\text{eff}}\). The errors in the experiment came from two major sources; the counting rate error and the uncertainties on the thickness of target. The maximum errors in the total mass attenuation coefficients were calculated from errors in incident (I\(_0\)) and transmitted (I) intensities and areal density (t) by using the propagation of error formula,

\[
\Delta \left( \frac{\mu}{\rho} \right) = \frac{1}{t} \left[ \left( \frac{\Delta I_0}{I_0} \right)^2 + \left( \frac{\Delta I}{I} \right)^2 + \left( \frac{\Delta t}{t} \right)^2 \right] \tag{11}
\]

where \(\Delta I_0\), \(\Delta I\), and \(\Delta t\) are the errors in the intensities \(I_0\), I, and thickness t of the sample, respectively. The error occurred in measured \((\mu/\rho)c\) is 2–3%.

**Conclusion**

We have studied the x-ray and gamma radiation shielding properties of Barium compounds such as \(\text{BaCl}_2\), \(\text{BaCO}_3\), \(\text{Ba(NO}_3)_2\) and \(\text{BaSO}_4\). The mass attenuation coefficient and linear attenuation coefficient for \(\text{BaCO}_3\) is higher than the \(\text{BaCl}_2\), \(\text{Ba(NO}_3)_2\) and \(\text{BaSO}_4\). HVL, TVL and mean free path are lower for \(\text{BaCO}_3\) than the \(\text{BaCl}_2\), \(\text{Ba(NO}_3)_2\) and \(\text{BaSO}_4\). Among the studied barium compounds, \(\text{BaCO}_3\) is best material for x-ray and gamma shielding.

**References**

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