A Study of Shielding Properties of X-ray and Gamma in Barium Compounds

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ABSTRACT

Background: Ionizing radiation is known to be harmful to human health. The shielding of ionizing radiation depends on the attenuation which can be achieved by three main rules, i.e. time, distance and absorbing material.

Materials and Methods: The mass attenuation coefficient, linear attenuation coefficient, Half Value Layer (HVL) and Tenth Value Layer (TVL) of X-rays (32 keV, 74 keV) and gamma rays (662 keV) are measured in Barium compounds.

Results and Discussion: The measured values agree well with the theory. The effective atomic numbers (Z_{eff}) and electron density (Ne) of Barium compounds have been computed in the wide energy region 1 keV to 100 GeV using an accurate database of photon-interaction cross sections and the WinXCom program.

Conclusion: The mass attenuation coefficient and linear attenuation coefficient for BaCO_3 is higher than the BaCl_2 , $\text{Ba}(\text{NO}_3)_2$ and BaSO_4 . HVL, TVL and mean free path are lower for BaCO_3 than the BaCl_2 , $\text{Ba}(\text{NO}_3)_2$ and BaSO_4 . Among the studied barium compounds, BaCO_3 is best material for x-ray and gamma shielding.

Keywords: X-ray, Gamma, Barium compounds

Original Research

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Introduction

Ionizing radiation is known to be harmful to human health. The shielding of ionizing radiation depends on the attenuation which can be achieved by three main rules, i.e. time, distance and absorbing material. Shielding is the most effective technique and it is required against X-rays and gamma rays, which are very penetrating [1, 2]. Gamma radiation is best absorbed by dense materials consisting of heavy atoms such as lead and barium [3]. Barium is a dense alkaline earth metal in Group IIA- (alkaline earth metals) of the periodic table that occurs in nature as a divalent cation in combination with other elements. Two commonly found forms of barium are barium sulfate and barium carbonate, often found as underground ore deposits.

Shielding aprons manufactured from environmentally-friendly barium sulfate as a potential substitute for the lead aprons. Barium sulfate shielding can be processed easily to any type, with the same shielding ability as lead. Barium sulfate has long been utilized for in-vivo tests by employing the difference in density between soft tissues in radiology, and it is one of the most economic environment-friendly shielding materials [4]. Hubbell-Seltzer [5] tabulated the mass attenuation co-efficient for elements of

atomic number $Z = 1$ to 92 and for energy range 1 keV-100 GeV. Berger- Hubbell [6] developed a XCOM program to calculate a mass attenuation coefficients and it is based on the Hubbell-Seltzer theoretical data. Later, Gerward et al. [7, 8] transformed XCOM program to the Windows platform under the name WinXCom.

It is necessary to search for new radiation shielding materials to replace more toxic and dense shielding materials such as lead. Barium is less toxic and cost effective material compared to lead. Due to the importance of Barium compounds in radiation shielding, the study of an interaction of gamma radiation in these compounds becomes important. Hence in the present work, we have studied radiation shielding properties of Barium compounds. The photon attenuation coefficients (linear and mass attenuation coefficients), effective atomic number and electron density are basic quantities required in determining the attenuation of X-rays and gamma photons in matter. The aim of the present work is to measure the X-ray (32 keV, 74 keV) and Gamma (662 keV) shielding parameters such as mass attenuation coefficient, and linear attenuation coefficient, Half Value Layer (HVL) and Tenth Value Layer (TVL) in Barium compounds such as BaCl_2 , BaCO_3 , $\text{Ba}(\text{NO}_3)_2$ and BaSO_4 . The effective atomic numbers (Z_{eff}) and electron density (N_e) of Barium have been computed in the wide region 1 keV to 100 GeV using an accurate database of photon-interaction cross sections and the WinX-Com program.

Materials and Methods

1. Present work

1) Theoretical calculation of effective atomic number and electron density

In the present work, the mass attenuation coefficients and photon interaction cross sections in the energy range from 1 keV to 100 GeV are generated using WinXCom [7, 8]. The total molecular cross section σ_m [milli barn], is computed from the following equation using the values of mass attenuation coefficients $[(\mu/\rho)_c]$

$$\sigma_m(E) = \left(\frac{1}{N} \right) \left(\frac{\mu(E)}{\rho} \right) \sum_i n_i A_i \quad (1)$$

where n_i is the number of atoms of i^{th} element in a given molecule, $(\mu/\rho)_c$ is the mass attenuation coefficient of biomolecule, N is the Avogadro's number, and A_i is the atomic weight of element i . The effective (average) atomic cross sec-

tion for a particular atom in the compound σ_a [milli barn] is estimated using the equation,

$$\sigma_a = \frac{\sigma_m}{\sum_i n_i} \quad (2)$$

The effective electronic cross section σ_e [milli barn] is computed from mass attenuation coefficient $(\mu/\rho)_i$ of i^{th} element in the given molecule using following equation,

$$\sigma_e = \left(\frac{1}{N} \right) \sum_i \left\{ \left(\frac{f_i A_i}{Z_i} \right) \left(\frac{\mu}{\rho} \right)_i \right\} = \frac{\sigma_a}{Z_{\text{eff}}} \quad (3)$$

where, f_i is the fractional abundance (a mass fraction of the i^{th} element in the molecule) and Z_i is the atomic number of the i^{th} element in a molecule. Finally the Z_{eff} is estimated as

$$Z_{\text{eff}} = \frac{\sigma_a}{\sigma_e} \quad (4)$$

The effective electron density, N_e , expressed in terms of number of electrons per unit mass is closely related to the effective atomic number. For an element, the electron density is given by $N_e = NZ/A$ [electrons/g]. This expression can be generalized for a compound,

$$N_e = \frac{N}{\sum_i n_i A_i} Z_{\text{eff}} \sum_i n_i \quad (5)$$

In the present work, N_e and Z_{eff} of Barium compounds are computed from Equations (4) and (5).

2) Measurement of mass attenuation coefficient of gamma and other derived gamma attenuation parameters

Transmission experiments with the narrow beam (good-geometry) setup were used for measuring the incident and transmitted intensities, and hence calculating the attenuation coefficient. The narrow geometry experimental setup used in the present measurement is as shown in Figure 1. We have used a NaI(Tl) crystal detector mounted on a photo-

Fig. 1. Schematic diagram of the Experimental Setup. S: Source position, T: Target sample, L: Lead shielding, D: Detector, PM: Photomultiplier.

multiplier tube housed in a lead chamber and a sophisticated PC based MCA for a detection purpose, gamma source such as ^{137}Cs (661.6 keV) and Barium compounds such as BaCl_2 , BaCO_3 , $\text{Ba}(\text{NO}_3)_2$ and BaSO_4 as target samples. The sample was directly attached to the opening of the lead shield where source is placed. The integral intensities, I_0 and I of the beam before and after passing through the sample are measured for sufficient time. $(\mu/\rho)_c$ of the sample is then estimated using the relation.

$$\left(\frac{\mu}{\rho}\right)_c = \left(\frac{I}{I_0}\right) \ln\left(\frac{I_0}{I}\right) \quad (6)$$

Where, t and ρ are the thickness and density of the sample respectively. Experimental values of N_e and Z_{eff} of Barium compounds are obtained by substituting the measured values of $(\mu/\rho)_c$ in Equations (1)-(5). Some of the earlier work-

ers measured the mass attenuation coefficient of photons using alternate to transmission experiments [9, 10]. The high-purity Germanium spectrometers having poor efficiency but high resolution, NaI scintillation detectors have a poor energy resolution but high efficiency [11]. For intensity measurements, efficiency of the detector should be high rather than resolution. Many previous [12-15] workers used the NaI scintillation detectors for the measurement of mass attenuation coefficients. Hence we have used the NaI scintillation detectors for the measurements.

Results and Discussion

The measured ^{137}Cs gamma spectrum in different thickness of target samples such as BaCl_2 , BaCO_3 , $\text{Ba}(\text{NO}_3)_2$ and BaSO_4 are shown in the Figures 2-5. The experimental mass

Fig. 2. The measured ^{137}Cs spectrum in different thickness of BaCl_2 sample.

Fig. 4. The measured ^{137}Cs spectrum in different thickness of $\text{Ba}(\text{NO}_3)_2$ sample.

Fig. 3. The measured ^{137}Cs spectrum in different thickness of $\text{Ba}(\text{CO}_3)_2$ sample.

Fig. 5. The measured ^{137}Cs spectrum in different thickness of BaSO_4 sample.

attenuation coefficients are compared with the theoretical values and it is shown in the Table 1. This comparison shows that the measured values agree well with the theoretical values.

The total linear attenuation coefficient (μ) can be evaluated by multiplying density of compounds to mass attenuation coefficients.

$$\mu = \left(\frac{\mu}{\rho} \right)_c \times \rho \quad (7)$$

The total linear attenuation coefficient (μ) is used in the calculation of HVL. HVL is the thickness of a shield or an absorber that reduces the radiation level by a factor of 2 that is to half the initial level and is calculated by the following equation

$$HVL = \frac{\ln 2}{\mu} = \frac{0.693}{\mu} \quad (8)$$

The total linear attenuation coefficient (μ) is also used in the calculation of TVL. It is the thickness of a shield required for attenuating a radiation beam to 10% of its radiation level and is computed by

$$TVL = \frac{\ln 10}{\mu} = \frac{2.303}{\mu} \quad (9)$$

The average distance between two successive interactions is called the relaxation length (λ). It is also called the photon mean free path which is determined by the equation:

$$\lambda = \frac{\int_0^\infty x \cdot \exp(-\mu x) dx}{\int_0^\infty \exp(-\mu x) dx} = \frac{1}{\mu} \quad (10)$$

The measured linear attenuation coefficients (cm^{-1}), HVL (in cm), TVL (in cm) and mean free path (in cm) are compared with the theoretical values. This comparison is as shown in the Tables 2-5. The theoretical mass attenuation coefficients for the given compounds are computed for wide energy range 1 keV-100 GeV. These values are used in the calculation of linear attenuation coefficients (cm^{-1}), HVL (in cm), TVL (in cm) and mean free path (in cm) for wide energy range 1 keV-100 GeV. Figures 6-8 show that the variation of HVL, TVL and mean free path for different barium compounds with energy. From these figures, it is found that for all compounds the values of HVL, TVL and mean free path re-

Table 1. Comparison of Measured Mass Attenuation Coefficients (cm^2/g) with Theoretical Values

Target compound	32 keV		74 keV		661.7 keV	
	Expt.	Theory	Expt.	Theory	Expt.	Theory
BaCl ₂	8.434±0.194	7.676	3.132±0.072	2.910	0.081±0.0019	0.0754
BaCO ₃	10.454±0.261	9.523	3.363±0.084	3.143	0.085±0.0021	0.0799
Ba(NO ₃) ₂	6.996±0.182	6.431	2.751±0.072	2.553	0.082±0.0022	0.0763
BaSO ₄	7.441±0.201	6.890	2.768±0.080	2.612	0.083±0.0024	0.0781

Table 2. Comparison of Measured Linear Attenuation Coefficients (cm^{-1}) with Theoretical Values

Target compound	32 keV		74 keV		661.7 keV	
	Expt.	Theory	Expt.	Theory	Expt.	Theory
BaCl ₂	32.557±0.814	29.629	12.091±0.302	11.233	0.314±0.008	0.291
BaCO ₃	44.849±1.166	40.854	14.426±0.375	13.483	0.363±0.009	0.343
BaNO ₃	31.483±0.866	28.940	12.381±0.340	11.489	0.367±0.010	0.343
BaSO ₄	24.107±0.686	22.324	8.970±0.255	8.463	0.268±0.008	0.253

Table 3. Comparison of HVL (cm) with Theoretical Values

Target compound	32 keV		74 keV		661.7 keV	
	Expt.	Theory	Expt.	Theory	Expt.	Theory
BaCl ₂	2.13E-02±5.32E-04	2.34E-02	5.73E-02±1.43E-03	6.17E-02	2.21E+00±5.51E-02	2.38E+00
BaCO ₃	1.55E-02±3.86E-04	1.70E-02	4.80E-02±1.20E-03	5.14E-02	1.91E+00±4.77E-02	2.02E+00
BaNO ₃	2.20E-02±5.50E-04	2.40E-02	5.60E-02±1.40E-03	6.03E-02	1.89E+00±4.72E-02	2.02E+00
BaSO ₄	2.88E-02±7.19E-04	3.10E-02	7.73E-02±1.93E-03	8.19E-02	2.58E+00±6.46E-02	2.74E+00

Table 4. Comparison of TVL (cm) with Theoretical Values

Target compound	32 keV		74 keV		661.7 keV	
	Expt.	Theory	Expt.	Theory	Expt.	Theory
BaCl ₂	7.07E-02 ± 1.77E-03	7.77E-02	1.90E-01 ± 4.76E-03	2.05E-01	7.33E+00 ± 1.83E-01	7.91E+00
BaCO ₃	5.13E-02 ± 1.28E-03	5.64E-02	1.60E-01 ± 3.99E-03	1.71E-01	6.34E+00 ± 1.59E-01	6.72E+00
BaNO ₃	7.31E-02 ± 1.83E-03	7.96E-02	1.86E-01 ± 4.65E-03	2.00E-01	6.27E+00 ± 1.57E-01	6.71E+00
BaSO ₄	9.55E-02 ± 2.39E-03	1.03E-01	2.57E-01 ± 6.42E-03	2.72E-01	8.59E+00 ± 2.15E-01	9.10E+00

Table 5. Comparison of Photon Mean Free Path (cm) with Theoretical Values

Target compound	32 keV		74 keV		661.7 keV	
	Expt.	Theory	Expt.	Theory	Expt.	Theory
BaCl ₂	3.07E-02 ± 7.68E-04	3.38E-02	8.27E-02 ± 2.07E-03	8.90E-02	3.18E+00 ± 7.96E-02	3.44E+00
BaCO ₃	2.23E-02 ± 5.58E-04	2.45E-02	6.93E-02 ± 1.73E-03	7.42E-02	2.76E+00 ± 6.89E-02	2.92E+00
BaNO ₃	3.18E-02 ± 7.94E-04	3.46E-02	8.08E-02 ± 2.02E-03	8.70E-02	2.72E+00 ± 6.81E-02	2.91E+00
BaSO ₄	4.15E-02 ± 1.04E-03	4.48E-02	1.12E-01 ± 2.79E-03	1.18E-01	3.73E+00 ± 9.32E-02	3.95E+00

Fig. 6. The variation of Half Value Layer (HVL) with energy for different barium compounds. **Fig. 8.** The variation of Tenth Value Layer means free path or relaxation length with energy for different barium compounds. **Fig. 7.** The variation of Tenth Value Layer (TVL) with energy for different barium compounds.

mains constant from 1 keV to 10 keV. These parameters increase sharply up to few MeV and thereafter shows decreasing trend. From these figures it is observed that the HVL, TVL and mean free path increases up to the E_{pe} and then decreases [16]. Here E_{pe} is the energy value at which the photo electric interaction coefficients matches with Compton interaction coefficients. The variation of these parameters with energy is due to dominance of photoelectric absorption in the lower end and dominance of pair production in the higher photon energy region. In the lower energy end photoelectric absorption is dominant, hence these values are minimum. As the energy of incident photon increases, Compton scattering overtakes the photoelectric absorption. It results multiple Compton scattering events which increases these values up to the E_{pe} and becomes maximum at E_{pe} . Thereafter

Fig. 9. The variation of Z_{eff} with energy for different barium compounds.

(above E_{pe}), pair production starts dominating (absorption process) which reduces these parameters to minimum value.

The variation of Z_{eff} with photon energy for total photon interactions is as shown in the figure 9 and this variation is because of dominance of different photon interactions. Z_{eff} increases and becomes maximum and decreases sharply in the energy region 0.002-0.35 MeV. The variation of Z_{eff} with energy is constant for 0.5 MeV to 2 MeV, From 2 MeV to 70 MeV increases and thereafter remains constant. The effective electron density shows similar trend as that of Z_{eff} . The errors in the experiment came from two major sources; the counting rate error and the uncertainties on the thickness of target. The maximum errors in the total mass attenuation coefficients were calculated from errors in incident (I_0) and transmitted (I) intensities and areal density (t) by using the propagation of error formula,

$$\Delta\left(\frac{\mu}{\rho}\right) = \frac{1}{t} \sqrt{\left(\frac{\Delta I_0}{I_0}\right)^2 + \left(\frac{\Delta I}{I}\right)^2 + \left(\ln \frac{\Delta I}{I}\right)^2 + \left(\frac{\Delta t}{t}\right)^2} \quad (11)$$

where ΔI_0 , ΔI , and Δt are the errors in the intensities I_0 , I , and thickness t of the sample, respectively. The error occurred in measured $(\mu/\rho)c$ is 2-3%.

Conclusion

We have studied the x-ray and gamma radiation shielding properties of Barium compounds such as BaCl_2 , BaCO_3 , $\text{Ba}(\text{NO}_3)_2$ and BaSO_4 . The mass attenuation coefficient and linear attenuation coefficient for BaCO_3 is higher than the

BaCl_2 , $\text{Ba}(\text{NO}_3)_2$ and BaSO_4 . HVL, TVL and mean free path are lower for BaCO_3 than the BaCl_2 , $\text{Ba}(\text{NO}_3)_2$ and BaSO_4 . Among the studied barium compounds, BaCO_3 is best material for X-ray and gamma shielding.

References

1. Stewart DC. Data for radioactive waste management and nuclear applications. New York, NY. John Wiley and Sons. 1985;89.
2. Chilton AB, Sheltie JK, Faw RE. Principles of radiation shielding. 2nd Ed. Englewood Cliffs, NJ. Prentice-Hall. 1984;105.
3. Dresner L. Principles of radiation protection engineering. 3rd Ed. New York, NY. McGraw-Hill. 1965;204.
4. Kim SC, Dong KR, Chung WK. Medical radiation shielding effect by composition of barium compounds. Ann. Nucl. Energy. 2012; 47;1-5.
5. Hubbell JH, Seltzer SM. Tables of x-ray mass attenuation coefficients and mass energy-absorption coefficients 1 keV to 20 MeV for elements Z=1 to 92 and 48 additional substances of dosimetric interest. NISTIR 5632. National Institute of Standards and Technology. 1995;56.
6. Berger MJ, Hubbell JH. XCOM: Photon cross sections on a personal computer. NBSIR 87-3597. National Bureau of Standards, U.S. Department of Commerce. 1987;28.
7. Gerward L, Guilbert N, Jensen KB, Levring H. lead phosphate glass containing boron and lithium oxides as a shielding material for neutron- and gamma-radiation. Rad. Phys.Chem. 2001;60: 23.
8. Gerward L, Guilbert N, Jensen KB, Levring H. WinXCom—a program for calculating X-ray attenuation coefficients. Rad. Phys.Chem. 2004;71:653.
9. Chantler CT, Tran CQ, Barnea Z, Paterson D, Cookson DJ, Balic DX. Measurement of the X-ray mass attenuation coefficient of copper using 8.85–20 keV synchrotron radiation. Phys. Rev. A. 2001;(6):062506.
10. Tran CQ, Chantler CT, Barnea Z. X-Ray Mass attenuation coefficient of silicon: Theory versus experiment, Phys. Rev. Lett. 2003; 90(25):257401.
11. Hossain I, Sharip N, Viswanathan KK. Efficiency and resolution of HPGe and NaI(Tl) detectors using gamma-ray spectroscopy. Sci. Res. Essays. 2012;7(1):86-89.
12. Goswami B, Chaudhuri N. Measurements of gamma-ray attenuation coefficients. Phys. Rev. A. 1937;7(6):1912.
13. Manjunatha HC. Influence of gamma irradiation on conductivity of $\text{YBa}_2\text{Cu}_3\text{O}_7$. Rad. Phys.Chem. 2005;113:24– 27.
14. Rudraswamy B, Dhananjaya N, Manjunatha HC. Measurement of absorbed dose rate of gamma radiation for lead compounds. Nucl. Instrum. Methods Phys. Res., Sect. A. 2010;619:171–173.
15. Manjunathaguru V, Umesh TK. Effective atomic numbers and

electron densities of some biologically important compounds containing H, C, N and O in the energy range 145–1330 keV. J. Phys. B: At. Mol. Opt. Phys. 2006;39(18):3969.

16. Manjunatha HC, Rudraswamy B. A study of thickness and penetration depth dependence of specific absorbed fraction of energy in bone. Ann. Nucl. Energy. 2011;38(10):2271-2282.